

1 Deuteration for High-Resolution Detection of Protons in Protein 2 Magic Angle Spinning (MAS) Solid-State NMR

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Cite This: <https://doi.org/10.1021/acs.chemrev.1c00681>



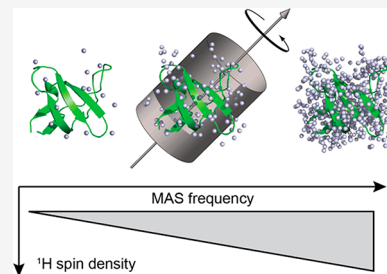
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4 **ABSTRACT:** Proton detection developed in the last 20 years as the method of choice to
5 study biomolecules in the solid state. In perdeuterated proteins, proton dipolar interactions
6 are strongly attenuated, which allows yielding of high-resolution proton spectra.
7 Perdeuteration and backsubstitution of exchangeable protons is essential if samples are
8 rotated with MAS rotation frequencies below 60 kHz. Protonated samples can be
9 investigated directly without spin dilution using proton detection methods in case the MAS
10 frequency exceeds 110 kHz. This review summarizes labeling strategies and the
11 spectroscopic methods to perform experiments that yield assignments, quantitative
12 information on structure, and dynamics using perdeuterated samples. Techniques for
13 solvent suppression, H/D exchange, and deuterium spectroscopy are discussed. Finally,
14 experimental and theoretical results that allow estimation of the sensitivity of proton detected experiments as a function of the MAS
15 frequency and the external B_0 field in a perdeuterated environment are compiled.



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1. INTRODUCTION

A The development of magic angle spinning (MAS) solid-state
B NMR has made tremendous progress in the past 20 years. 45
C Improved experimental performance was triggered by the
D development of ever faster spinning probes.^{1,2} Increased 46
E rotation frequencies result in a more efficient averaging of
F proton dipolar couplings, which in 2016 culminated in the
G determination of the first de novo protein structure which is
H based on proton detection by MAS solid-state NMR.³ Initially, 47
I multipulse homonuclear decoupling sequences such as
J CRAMPS,^{4–6} wPMLG,^{7,8} and wDUMBO⁹ have been employed 48
K to suppress proton homonuclear interactions in either the direct
L or indirect dimension. In CRAMPS type sequences, the signal is
M detected stroboscopically by including windows in the multiple
N pulse sequence. As such, the sequences are restricted to
O relatively low MAS rotation frequencies (<13 kHz). Windowed
P detection schemes come along with reduced sensitivity as large
Q filter band widths have to be employed. 49
R

S Alternatively, strong proton-proton dipolar couplings can be
T suppressed chemically by deuteration. There, protons are
U replaced with deuterons at all nonexchangeable positions. 50
V After protein expression, deuterons at exchangeable positions
W are replaced with protons. This approach was originally 51
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Z 54
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Special Issue: Biomolecular NMR Spectroscopy

Received: August 2, 2021

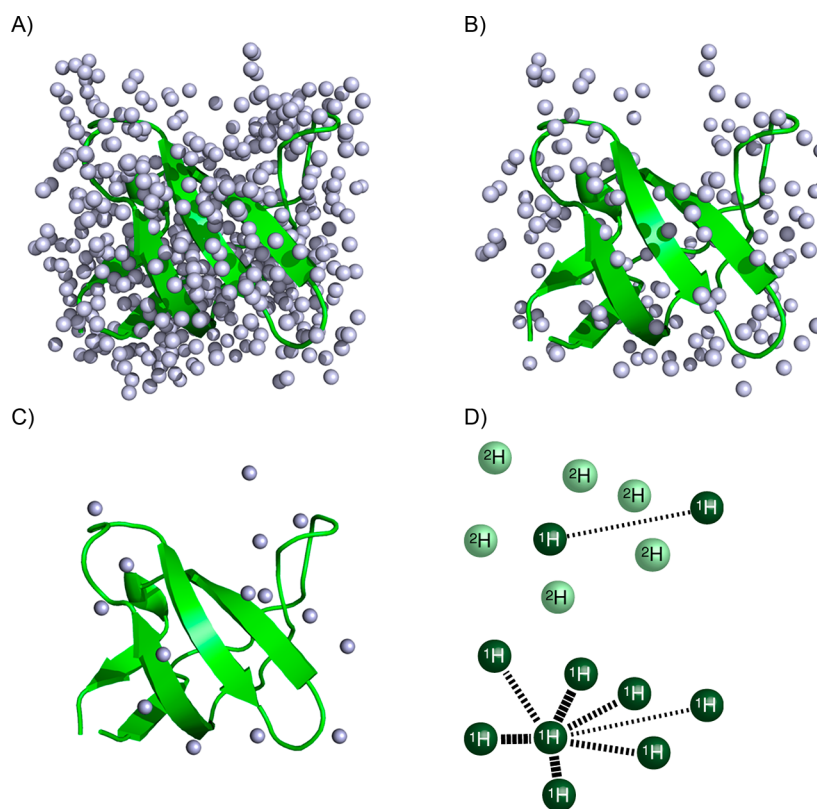


Figure 1. α -Spectrin SH3 domain at different levels of deuteration (PDB 1U06).¹⁰ In comparison to the protonated sample (A), the number of proton dipolar interactions is already significantly reduced in the perdeuterated sample that is recrystallized from a buffer that contains 100% H₂O (B). If D₂O is admixed into the crystallization buffer, the proton spin density is further diminished (C). The attenuated proton dipolar network allows to detect protons in the solid-state with high resolution (D). Reproduced with permission from ref 11. Copyright 2012 Elsevier.

introduced for solution NMR.^{12–14} In the solid state,
deuteration was first employed to study small molecules,^{15–17}
then later extended to the investigation of peptides^{18–20} and
proteins.^{21–24} Deuterons have a 6.5-fold smaller magnetogyric
ratio compared to protons. This facilitates homonuclear and
heteronuclear decoupling. The dilution of the proton spin
system by deuteration is illustrated in Figure 1.

2. NARROW PROTON RESONANCES BY SPIN DILUTION

2.1. Amide Protons

Spin dilution at exchangeable sites yields high resolution amide
proton spectra in perdeuterated proteins at MAS
rotation frequencies on the order of 10–20 kHz.²⁵ The
experimental proton line widths are on the order of 17–35
Hz, which is equivalent to 0.028–0.058 ppm at 600 MHz. Early
examples and applications of this labeling strategy involve the
investigation of the membrane protein DsbB reconstituted in a
lipid environment.²⁶ In this work, the resonances of DsbB were
assigned using ¹H-detected 3D NMR experiments. Lakomek et
al. investigated and compared solid- and solution-state lipid
nanodisc preparations in which the outer-membrane protein
OmpX was reconstituted.²⁷ Furthermore, proton detection
experiments have been employed to probe the PET ligand
binding site in Alzheimer's disease A β 40 fibril.²⁸ For hepatitis B
virus capsids, it has been shown that back-substitution of
exchangeable deuterons with amide protons improves the
proton resolution at 100 kHz MAS by a factor of 1.5
in comparison to a protonated sample.²⁹ The average proton line

width decreases from 170 to 110 Hz, which corresponds to a
decreased from 0.20 to 0.13 ppm at 850 MHz. However, because
of solvent inaccessibility, 15% of all amides are absent in the
spectra of the deuterated protein because exchangeable protons
cannot easily be back-substituted in the reconstitution process.

2.2. Methyl Protons

Similarly, methyl protons can be observed at high resolution in
perdeuterated peptides and proteins.³⁰ The media that are
employed to grow bacteria and overproduce the protein of
interest contain glucose that is labeled only with an efficiency of
~97–98% with deuterium. Thus, there is a 6–9% probability to
find a proton in a methyl group. At 600 MHz (14.1 T) and using
a MAS rotation frequency of 22 kHz, an experimental proton
(carbon) line width of 20–25 Hz (5–8 Hz) was observed. Use
of specific amino acid precursors such as pyruvate³¹ and α -
ketoisovalerate^{32–34} during protein biosynthesis allows selective
labeling of methyl groups and increasing of the sensitivity
further. For liquid-state NMR, this labeling strategy was
pioneered by Kay and co-workers.³⁵ Even higher performance
is obtained by making use of stereospecific precursors which
yield enrichment in only one of the two pro-chiral methyl groups
in valine, leucine, or isoleucine³⁶ as well as for methionine,
threonine, and alanine.³⁷ In the solid state, the intramethyl
proton dipolar couplings can potentially impact resolution.
Therefore, CHD₂ isotopomers should be incorporated into the
protein.³⁸ In addition, care has to be taken to decouple
deuterium scalar couplings while evolving carbon chemical
shifts. The ²H,¹³C scalar coupling is on the order of 20 Hz. As
deuterium is a spin-1 nucleus, scalar couplings contribute
significantly to the achievable ¹³C line width.

2.3. Other Aliphatic Sites

122 Randomly protonated (RAP) samples are obtained by growing
123 bacteria in a medium containing ^2H , ^{13}C labeled glucose in the
124 presence of varying amounts of H_2O (5–20%).³⁹ Using this
125 procedure, samples are obtained that are randomly protonated
126 in the side chain.

127 A M9 medium that is supplemented with 5% [15%] H_2O
128 yields approximately 7% [17%] protons in the $\text{C}\alpha$ position.
129 Under these conditions, a labeling efficiency of 3% [14%] and
130 1% [12%], respectively, is observed for methylene and methyl
131 groups. Side chain resonances are assigned in HCCH type
132 experiments in a straightforward manner.^{40,41} If faster MAS
133 frequencies are available, larger amounts of H_2O in the M9
134 medium are tolerable. In addition to line broadening that is
135 induced by proton–proton dipolar couplings within the methyl
136 group, mixtures of isotopomers result in an apparent broadening
137 of the resonances and in a deterioration of the spectral
138 resolution.

139 An admixture of 15 and 25% H_2O into the M9 medium yields
140 high sensitivity and resolution in the MAS frequency regime of
141 50 kHz. For these kinds of samples, an average ^1H line width of
142 40–50 Hz is observed.⁴² ^1H , ^{13}C correlation spectra of RAP
143 samples yield almost comparable intensities in comparison to
144 samples in which α -ketoisovalerate was used as a precursor.⁴³
145 RAP labeled samples are readily prepared and allow obtaining of
146 methyl spectra from all methyl containing amino acids without
147 the need to use specific amino acid precursors.³⁶ The RAP
148 labeling scheme has been successfully applied to a fibril sample
149 of the HET-s (218–289) prion protein.⁴⁴ An approach coined
150 inverse fractional deuteration employs protonated glucose and
151 D_2O to achieve a similar effect.^{45,46} Alternatively, protonated
152 amino acids are added to the growth medium that consists of
153 deuterated glucose and D_2O to selectively introduce protons
154 into protein.⁴⁷ Use of transamination enzymes allows to yield
155 quantitative protonation at the backbone $\text{H}\alpha$ position.⁴⁸ In the
156 approach, the keto acid is transformed by transaminases into the
157 respective amino acid. The proton that is bound to the $\text{C}\alpha$
158 carbon originates from the solvent. Racemization is avoided
159 because the correct L-amino acids are generated enzymatically.

2.4. Exchangeable Side Chain Protons

160 Hydroxyl protons exchange rapidly with the solvent. Due the
161 fact that they experience only weak couplings with their
162 environment, hydroxyl protons are difficult to assign in solution
163 NMR. By making use of dipolar couplings, magnetization can be
164 transferred approximately 50 \times faster in the solid state. In a long-
165 range ^1H , ^{13}C cross-polarization experiment, all threonine
166 hydroxyl protons could be assigned in the α -spectrum SH3
167 domain by correlating the OH proton and the $\text{C}\alpha/\text{C}\beta$
168 resonances.⁴⁹ An additional delay during which magnetization
169 is stored longitudinally yields the exchange properties of the
170 respective hydroxyl protons. Similarly, collective exchange
171 processes in the active-site proton cage in bacteriorhodopsin
172 could be characterized.⁵⁰ Furthermore, backbone hydrogen
173 bonds that stabilize protein secondary structure elements in
174 solid proteins could be identified.⁵¹ The same experiment
175 applied to side chain resonances showed that salt bridges
176 stabilize amyloid quaternary structure.⁵² By using a similar
177 strategy, the hydrogen bonding topology of histidines in the
178 influenza M2 channel has been identified.³⁴ More recently, the
179 hydrogen bonding topology could be characterized by
180 quantifying the scalar coupling across a hydrogen bond.⁵³

These experiments are carried out with a protonated sample that
is rotated with a frequency of 100 kHz. 181 182

3. SOLVENT SUPPRESSION

Biological solid-state samples contain a fair amount of aqueous
buffer to keep samples hydrated. In ^1H detected experiments,
water suppression thus becomes an issue. It was shown that
pulsed field gradients can be employed to attenuate the water
resonance.²¹ Alternatively, a constant-time-like indirect evolu-
tion period in combination with proton decoupling allows
efficient cancellation of unwanted solvent magnetization.²⁴ The
MISSISSIPPI sequence⁵⁴ combines these two approaches and is
nowadays the most widespread employed pulse sequence
element to achieve efficient water suppression. Gradients are
not essential for the MISSISSIPPI sequence but improve water
suppression significantly. Care has to be taken to avoid an
attenuation of the resonances of the protein by a transfer of
saturation from the water pool. Therefore, water suppression
should be combined with water flip-back pulse schemes which
return the water magnetization back along the z-axis prior to
detection.⁵⁵ Water flip-back pulses facilitate spin recovery and
enable faster recycle delays. At the same time, a potential
saturation transfer from water to the protein resonances is
avoided. 183 184 185 186 187 188 189 190 191 192 193 194 195 196 197 198 199 200 201 202

4. PARAMAGNETIC DOPING, SOLVENT ACCESSIBILITY, AND MOLECULAR INTERFACES

Sensitivity can be improved by reducing the recycle delay of the
experiment making use of complexed paramagnetic ions, such
as, e.g., Cu-EDTA.^{56–58} This way, the recycle delay in the
experiment can be decreased by up to 15-fold. At the same time,
the proton line shape is largely unaffected. Perdeuterated
samples do not require high-power proton decoupling. Thus,
long acquisition times and short recycle delays on the order of
0.3 s are possible without compromising the sample quality. The
paramagnetic relaxation agent in solution can be employed to
furthermore access solvent-PREs. Solvent-PREs were originally
introduced in solution NMR to aid the protein structure
determination process⁵⁹ and were exploited in the solid state to
characterize protein–protein interfaces.⁶⁰ The solvent acces-
sible surface can be characterized by quantification of the
difference $^1\text{H}-T_1$ relaxation rate that is obtained for samples with
and without complexed paramagnetic ions.⁶¹ Alternatively,
molecular interfaces can be probed by comparing the spectra of a
perdeuterated, ^{15}N labeled protein in the presence of a
deuterated versus a protonated ligand. The differential broad-
ening of amide resonances due to dipolar interactions with the
ligand were exploited to identify the interface between a box C/
D RNA and the binding L7Ae protein.⁶² Similarly, the ligand
binding interface can be determined by making use of
magnetization transfer from a protonated ligand to the carbon
resonances of a perdeuterated protein, as shown for Congo Red
interacting with HET-s(218–289) prion fibrils.⁶³ Similarly, the
interface between CAP-Gly and microtubules could be
determined. This was achieved by making use of dREDOR
filtered experiments.⁶⁴ 204 205 206 207 208 209 210 211 212 213 214 215 216 217 218 219 220 221 222 223 224 225 226 227 228 229 230 231 232

In case the concentration of protons at exchangeable sites
cannot be adjusted prior to the NMR experiment by refolding
the protein in an appropriate buffer, perdeuteration can be an
issue because certain amides in structured regions might be
strongly protected from exchange. Back-exchange of deuterons
with protons then becomes impossible. On the other hand, the
233 234 235 236 237 238

239 absence of certain amide resonances allows identification of the
240 solvent-exposed parts of the protein, as has been shown for the
241 seven-helical integral membrane proton pump proteorhodop-
242 sin.⁶⁵ In addition to residues located at the membrane interface,
243 residues are also observed that are located in helix G, suggesting
244 that the F–G loop may have a high mobility and transiently
245 expose a hydrophilic cavity in the extracellular half of the protein
246 which might be part of a proton-conducting pathway. Similarly,
247 the spectra of fully protonated and perdeuterated proteorho-
248 dopsin were compared to yield the assignment of the nonsolvent
249 accessible core of the protein.⁶⁶

250 Direct interactions between the solvent and the human
251 voltage dependent anion channel, and the alkane transporter
252 AlkL, respectively, are probed by recording a 3D HhNH
253 correlation spectrum in which proton–proton mixing is
254 achieved via spin diffusion.⁶⁷ Similar experiments were carried
255 out for the potassium channel NaK2K.⁶⁸ It was shown that the
256 conduction pathway of the potassium channel is free of water
257 under physiological conditions. Lange and co-workers have
258 furthermore investigated the structure and dynamics of the
259 rhomboid protease GlpG reconstituted in liposomes and have
260 confirmed the presence of water molecules in the catalytic
261 cavity.⁶⁹ Dynamics measurements revealed a dynamic hotspot of
262 GlpG that is important for gating. Alternatively, the binding
263 interface can be probed by analyzing the differential dynamics in
264 the presence and absence of a ligand. This has been shown for
265 protein G in complex with a full-length human immunoglobu-
266 lin.^{70,71} In addition to chemical shift perturbations, R_1 and $R_{1\rho}$
267 relaxations rates suggest an increased amount of slow motions
268 with a time scale in the >500 ns range in the complex.

269 In a very favorable case, a proton detected experiment allowed
270 to detect a structured water molecule in the conductance
271 domain of the homotetrameric influenza A M2 membrane
272 protein.⁷² The bound water molecule was observed at a proton
273 chemical shift of 11 ppm and was found to be hydrogen bonded
274 to H37- δ 1 in the channel.

5. LIMITS OF RESOLUTION IN THE SOLID STATE

275 In the solid state, the NMR line width is dependent on several
276 factors. Inhomogeneous broadening⁷³ results in a distribution of
277 chemical shifts, for example, in a sample which is structurally
278 heterogeneous. By contrast, homogeneous broadening is a
279 consequence of dynamics and thus relaxation. In the following,
280 the factors that affect resolution at the upper limit where sample
281 heterogeneity is not an issue shall be discussed.

282 The decoupling performance and the accuracy of the magic
283 angle have a major impact on the resonance line width in solid-
284 state NMR. For perdeuterated samples, however, the decoupling
285 performance is not compromising the spectral quality because
286 long acquisition times can be employed without damaging the
287 sample. Mis-adjustments of the sample rotation angle results in a
288 broadening of the resonance.^{74,75} The spinning angle can be
289 optimized by monitoring the H–N J -coupling spin echo signal
290 in the protein sample under investigation.⁷⁶ Spin-state selective
291 experiments that select a particular multiplet component, for
292 which the ^1H – ^{15}N dipole–dipole interaction and the ^{15}N CSA
293 mutually cancel each other at high field due to interference of
294 anisotropic interactions allow furthermore to overcome magic
295 angle mis-settings.⁷⁷ The quality of the shim is another factor
296 that potentially contributes to line broadening. However, the
297 residual proton line width of the water resonance is less than 1
298 Hz in a typical microcrystalline protein sample, suggesting that
299 the homogeneity of the magnetic field does not limit the

resolution of the ^1H detected experiments. On the other hand, 300
the proton dipolar network directly impacts the resolution in a 301
 ^1H detected experiment. At a certain degree of deuteration and 302
MAS frequency, however, the spectral quality cannot be 303
improved further,^{25,78–81} suggesting that proton dipolar 304
interactions are not the critical factor that determines the 305
resolution of the proton spectrum. Protein dynamics and thus 306
incoherent effects have a further influence on the proton line 307
width. The experimental relaxation data can be employed to 308
calculate the resonance line width that results from dynamics.⁸² 309
It is found, however, that dynamics contributes less than 2 Hz to 310
line width for most residues in the α -spectrin SH3 domain. 311

Another factor that potentially affects line width is the 312
anisotropic bulk magnetic susceptibility (ABMS).^{1,83} In a 313
powder sample, the individual crystallites are nonhomoge- 314
neously distributed. Crystallites and mother liquor have a 315
distinct bulk magnetic susceptibility (BMS). The induced 316
magnetic polarization from the external magnetic field produces 317
a screening field that results in an inhomogeneously broadened 318
resonance line. The crystallites appear as magnetic dipoles in the 319
liquid. If the microcrystallites would be spherical, the BMS 320
would be isotropic and could be fully spun out by MAS. 321
However, given the fact that microcrystals, fibrils, or lipid 322
reconstituted membrane proteins are not spherical, the strength 323
of the induced dipolar fields depends on their orientation with 324
respect to the external magnetic field. The ABMS Hamiltonian 325
has a similar dependence on the nuclear spin part as the isotropic 326
chemical shift.¹ Application of a Carr–Purcell–Meiboom–Gill 327
(CPMG) train of pulses yields a significantly extended 328
coherence lifetime.^{84–86} A significantly extended FID is 329
obtained if CPMG pulses are applied in either the indirect or 330
direct evolution period in a ^1H , ^{15}N correlation experiment.⁸⁷ 331
This suggests that ABMS is responsible for the finite line width 332
in the solid-state. 333

6. ASSIGNMENT EXPERIMENTS

Assignments of perdeuterated and back-exchanged proteins in 334
rotating solids is nowadays routinely achieved using the suite of 335
experiments suggested by Barbet-Massin et al.⁸⁸ These experi- 336
ments are employed as well now for fully protonated samples at 337
>110 kHz MAS.^{3,89} Alternative strategies are suggested by 338
Lange,⁹⁰ Meier,⁹¹ and co-workers. 339

Coherences in the solid state are long-lived enough to enable 340
 J -coupling based assignment experiments for perdeuterated 341
samples such as 3D HNCA, HNCOC, and HNCACB.^{92,93} Penzel 342
et al. have compared scalar and dipolar coupling based transfer 343
sequences.⁹¹ Even scalar couplings across hydrogen bonds in 344
proteins can be quantified reliably.^{53,94} This experiment is 345
intrinsically very insensitive due to the necessary long de- and 346
refocusing delays. In contrast to solution NMR, however, the 347
achievable resolution and line width in the solid state are not 348
determined by nuclear relaxation implied by molecular 349
tumbling. Therefore, small scalar couplings are in principle 350
accessible in the solid state. 351

Barbet-Massin and co-worker mostly employ scalar transfer 352
elements to mediate homonuclear magnetization transfer. 353
Alternatively, band-selective homonuclear cross-polarization 354
(BSH-CP) can be employed to yield transfer of magnetization 355
between the carbonyl and $C\alpha$ carbon atoms.^{95,96} The experi- 356
ment works well for high magnetic fields and MAS frequencies 357
on the order of 20 kHz. Under these conditions, the ^{13}C , $^{13}\text{C}\alpha$ 358
isotropic chemical shift difference is greater than the MAS 359
frequency. Recoupling is obtained when the sum of the effective 360

361 rf fields [$\omega_1^{\text{eff}}(^{13}\text{CO}) + \omega_1^{\text{eff}}(^{13}\text{C}\alpha)$] matches $2 \times \omega_R$. This
362 method can directly be employed in assignment experiments
363 such as a 3D-(H)CO(CA)NH experiments to obtain inter-
364 residual connectivities in the protein backbone. The assignment
365 process can be supplemented with amide–amide connectivities
366 obtained from (H)N(CACO)NH or (H)N(COCA)NH type
367 3D^{97,98} and 4D⁹⁹ experiments. The resulting connectivities are
368 diagonal-free and can be employed for automatic backbone
369 assignment, e.g., in UNIO-MATCH.^{100,101}

370 Combination of direct ^{13}C Boltzmann polarization and
371 magnetization transferred from directly bonded or remote
372 protons can increase the sensitivity of the experiment.^{102,103}
373 Carbon mixing over multiple bonds using TOBSY¹⁰⁴ or
374 MOCCA¹⁰⁵ sequences enables access to virtually all side chain
375 carbon chemical shifts and can be detected either on the
376 methyl³⁰ or amide proton.^{102,106}

377 To speed up the sampling process, GFT and APSY have been
378 introduced to biological solid-state NMR. GFT projection
379 experiments are based on synchronous evolution of multiple
380 nuclei in a single dimension. The convoluted spectra are added
381 and subtracted to yield the reconstructed pure chemical shift
382 spectra. GFT projection experiments have been recorded from
383 4D/3D (HA)CANCOX and 3D/2D (HACA)NCOX
384 experiments.¹⁰⁷ Automated projection spectroscopy (APSY)
385 was originally introduced by Hiller and Wüthrich for solution
386 NMR^{108,109} and recently applied in the solid state.^{110,111} The 5D
387 peak lists that are reconstructed from a number of 2D
388 projections of a 5D-(H)NCOCANH experiment are a next
389 step toward an automated assignment of the protein back-
390 bone.¹¹¹ The efficiency of the pulse scheme is ensured by long
391 $T_{1\rho}$ and T_2' coherence lifetimes in the solid state. Additionally,
392 pairs of 4D experiments that correlate HN_i/N_i with CA_i/CO_i or
393 $\text{CA}_{i-1}/\text{CO}_{i-1}$ allow to obtain sequential connections with high
394 confidence.^{112,113}

7. DEUTERIUM SPECTROSCOPY

395 In the perdeuteration approach described here, deuterons are
396 primarily introduced into the sample to attenuate the ^1H dipolar
397 network to achieve proton line narrowing. On the other hand,
398 deuteron can be employed actively for spectroscopy in the solid
399 state. Given the fact that overall tumbling is absent in
400 immobilized systems, high-resolution ^2H , ^{13}C correlation experi-
401 ments can be obtained.^{75,114} ^2H double quantum (DQ)
402 coherences yield an improved spectral resolution.¹¹⁵ Sensitivity
403 is not compromised by missetting of the rotation axis from the
404 magic angle. At the same time, coherences are less sensitive to
405 dynamics that potentially interferes with refocusing of the
406 anisotropic interaction by MAS. DQ spectroscopy effectively
407 doubles the spectral resolution because ^2H DQ coherences
408 precess 2-fold faster in comparison to single quantum (SQ)
409 coherences.⁷⁵

410 Deuterium is the paradigmatic nucleus to study molecular
411 motion in solids.^{116,117} Deuterium has been employed in the
412 past to investigate synthetic polymer,¹¹⁸ lipid,^{119,120} and protein
413 dynamics^{121,122} of static samples. Using small increments in the
414 indirect evolution period allows reintroduction of the deuterium
415 Pake pattern under slow MAS and to retrieve dynamic
416 information.^{123–127} ^2H , ^2H spin diffusion potentially compro-
417 mises the deuterium spinning sideband pattern in ^2H , ^{13}C
418 correlation experiments and the extracted anisotropy param-
419 eters.^{128,129} Similarly, deuterium spectroscopy has been used to
420 study the dynamic properties of carbohydrates and bacterial
421 cellulose.¹³⁰

The ^2H RF field strength determines the ^2H , ^{13}C magnet- 422
ization transfer efficiency and should exceed >80 kHz for a rigid 423
deuteron. This power can so far only be achieved in MAS probes 424
that operate in double channel mode. Use of optimal control 425
(OC) derived pulse sequences allows this problem to be 426
partially overcome.¹³¹ Promising results were obtained with the 427
RESPIRATION scheme, where rf amplitudes were restricted to 428
a maximum of 50 kHz (with an average rf field of 11 kHz). 429
However, more work is needed to design pulse sequences that 430
use even lower rf fields and are able to access backbone $\text{C}\alpha$ or 431
other methine chemical groups. To fully exploit the potential of 432
deuterium in multidimension correlation experiments, a four- 433
channel MAS solid-state NMR probe is needed that is equipped 434
with a high power deuterium channel and a sensitive magic angle 435
adjustment, so that deuterium spectroscopy and proton 436
detection can be concatenated and combined with amide 437
correlation experiments.¹³² 438

Deuterons can serve as magnetization entry point for any 439
multidimensional NMR experiment. Combination of direct ^{13}C 440
Boltzmann polarization and magnetization transferred from 441
directly bonded or remote deuterons can increase the sensitivity 442
of the experiment.¹³³ 443

8. H/D EXCHANGE EXPERIMENTS

Different H/D exchange experiments that make use of 444
perdeuterated samples have been suggested recently. Lopez 445
del Amo et al.¹³⁴ propose an experiment which makes use of the 446
deuterium isotope effect on the amide nitrogen chemical shift 447
for perdeuterated protein samples that are recrystallized in a 448
buffer that contains a significant amount of D_2O . Exchange is 449
probed by direct ^{15}N detection or by concatenating two 450
exchange mixing times in combination with proton detection 451
and ^{15}N – ^1H / ^{15}N – ^2H isotope filtering elements. The experi- 452
ment is sensitive to H/D exchange times of up to 30 s, which 453
reflects the typical amide nitrogen relaxation time in a protein. 454
Deuterium back-exchange experiments at labile sites were 455
carried out for a perdeuterated, fully proton back-exchanged 456
Escherichia coli type I pili sample to investigate solvent 457
accessibility.¹³⁵ For that purpose, a preassembled reprotoneated 458
pili sample was washed several times with nondenaturing 100% 459
 D_2O buffer to remove all solvent accessible labile protons. It was 460
found that amide protons protected from H/D exchange are 461
mostly found in the core of the monomeric subunits of the pili, 462
while high H/D exchange rates are observed in lateral and axial 463
intermolecular interfaces. The HDX data thus allows to get a 464
better understanding of the stability and mechanical properties 465
of the pili. 466

Hydrogen exchange affects the ^{15}N longitudinal relaxation 467
time of amides in a protein. Comparison of ^{15}N - R_1 values 468
measured for samples that are prepared with different amounts 469
of D_2O in the buffer allows quantification of the exchange rate. 470
This approach is coined Relax-EXSY and was employed to 471
determine amino acid specific exchange rates in the type III 472
secretion system needle protein.^{136,137} The experiments 473
described above are equilibrium experiments. By contrast, 474
Grohe and co-worker have suggested a nonequilibrium 475
experiment to probe backbone amide hydrogen–deuterium 476
exchange rates in the solid state.¹³⁸ In the experiment, the 477
supernatant buffer of a 100% back-substituted microcrystalline 478
protein sample was replaced with a 80% deuterated buffer before 479
the sample was packed the into a 1.3 mm MAS rotor. Peak 480
intensities were monitored for 2 weeks to identify H/D 481
exchange in real time. 482

9. QUANTIFICATION OF DISTANCES

483 In the presence of strong dipolar couplings, the evolution of
484 weak interactions is attenuated. This effect is referred to a
485 dipolar truncation.^{139,140} As a consequence, it is difficult to
486 obtain long-range distance information in uniformly isotopically
487 labeled samples. In carbon-detected experiments using proto-
488 nated protein samples, dipolar truncation effects are avoided by
489 preparing carbon spin dilute samples. These samples can be
490 obtained by overexpressing proteins using a medium that
491 contains, e.g., either [1,3]-¹³C-glycerol or [2]-¹³C-glycerol as the
492 only source for carbon atoms.^{141,142} In fully isotopically labeled
493 samples, third-spin assisted recoupling mechanism can poten-
494 tially facilitate homonuclear long-range magnetization transfer
495 via the proton spin reservoir. This effect is exploited in
496 PAR^{143–145} and PAIN^{146,147} experiments. In perdeuterated
497 proteins, the proton spin bath is naturally dilute. Only
498 exchangeable sites are protonated. This labeling scheme thus
499 facilitates the quantification of long-range proton distance
500 restraints.^{18,22,23} C7¹⁴⁸ derived ¹H,¹H dipolar recoupling
501 experiments work well for microcrystalline peptides¹⁸ but are
502 not very efficient for microcrystalline protein samples,²²
503 presumably due to the presence of tightly bound water
504 molecules which perturb the idealized 2-spin system. Spin
505 diffusion,¹⁴⁹ RFDR,¹⁵⁰ or DREAM,¹⁵¹ like mixing sequences,
506 however, are suitable to recouple proton–proton dipolar
507 interactions at least at moderate rotation frequen-
508 cies.^{22,33,152–154} Even though proton RFDR and DREAM
509 have been successfully used at 100 kHz MAS,^{89,155} novel
510 recoupling schemes will be needed to increase the efficiency of
511 the magnetization transfer at even faster MAS frequencies.¹⁵⁶
512 In addition, selective labeling of methyl groups yields proton
513 spin dilution and access to methyl–methyl and methyl–amide
514 distance restraints.^{152,154,157,158}
515 For methyls, care has to be taken to label samples with CHD₂
516 isotopomers in case experiments are carried out with MAS
517 rotation frequencies below 50 kHz because dipolar interactions
518 are otherwise not sufficiently removed by MAS.¹⁵⁹ RAP labeling
519 yields randomly protonated samples, and this way allows to
520 obtain distance restraints for other aliphatic sites.³⁹ Beyond
521 methyls, selective labeling of amino acids is possible via the SAIL
522 approach.^{160,161}
523 Using proton–proton distance restraints, the structure of the
524 dimeric (2 × 16 kDa) Zn(II)-loaded superoxide dismutase
525 could be determined by solid-state NMR experiments.¹⁶² In
526 these experiments, RFDR¹⁶³ is employed to recouple proton
527 spins.
528 The NMR structure of an influenza A M2 variant, S31N, that
529 confers drug resistance, was obtained using selectively methyl
530 protonated protein samples.³⁴ There, methyl–methyl distance
531 restraints were obtained from a 4D proton detected ¹H–¹H
532 RFDR experiment. The experimental setup is similar to an
533 earlier implementation in which DREAM has been used for
534 proton–proton mixing.³³ Proton chemical shift information
535 from histidine side chains allowed to conclude further on the
536 mechanism of H⁺ transport which is critical for viral replication.
537 3D (H)NHH and (H)N(HH)NH correlation spectra of a
538 deuterated OmpG sample in which the exchangeable sites
539 contained protons close to 100% allowed to identify amide–
540 amide through-space connectivities and to determine the β-
541 sheet topology in the β-barrel membrane protein.¹⁶⁴
542 Exact eRFDR distance restraints are obtained in 3D ¹⁵N-
543 edited RFDR spectra (H-RFDR-hNH) experiments.¹⁶⁵ Quanti-

fication of the magnetization buildup in the initial-rate regime in
544 combination with a spin-diffusion correction is required to yield
545 exact distances. The exact distance restraints facilitate the
546 convergence of the structure calculation process and result in a
547 better defined structure.

Long-range distance restraints (up to ~2 nm) can be obtained
549 furthermore by measuring PREs. For superoxide dismutase,
550 PRE restraints were obtained from ¹⁵N-R₁ relaxation difference
551 rates measured using a Cu²⁺, Zn²⁺, and a Cu⁺,Zn²⁺ complexed
552 protein sample.¹⁵³ In combination with ¹H,¹H distance
553 restraints, PRE restraints were shown to be extremely valuable
554 to determine the structure of the protein in the solid state.
555 Furthermore, pseudocontact shifts can be employed as well in
556 the solid state to aid the structure determination process for
557 metalloproteins.^{166,167}

In case no natural metal binding sites are available, the protein
559 under investigation, paramagnetic centers can be engineered
560 into a protein side chain using, e.g., a suitable Cu(II)-affinity
561 tag.¹⁶⁸ Backbone amide ¹⁵N longitudinal and ¹H transverse
562 relaxation rates for the paramagnetic and diamagnetic samples
563 are quantified in proton-detected experiments to determine
564 distances between the paramagnetic center and the respective
565 amide sites in the protein. The sparse PRE restraints allow
566 calculation of highly accurate protein structures.^{169,170}

Using Gd³⁺ paramagnetic ions bound to molecular cages
568 chelated to the lipid headgroups, it was possible to determine the
569 orientation of the D1 domain of the membrane-associated
570 nonstructural protein 5A (NSSA) of hepatitis C virus with
571 respect to the lipid membrane.¹⁷¹

10. QUANTIFICATION OF DYNAMICS

NMR experiments allow obtaining access to quantitative
573 motional amplitudes and rates. Different solid-state NMR
574 experiments are sensitive to dynamics occurring on different
575 time scales: relaxation experiments yield information on fast to
576 intermediate time scale (ps–μs) motion. Exchange experiments
577 are sensitive to slow motional processes (>1 ms) when the
578 chemical shifts of resonances for different conformers are
579 resolved. Motion faster than 1 μs averages the anisotropic
580 nuclear spin interactions. The amplitude of the motional process
581 can thus be estimated from experiments that reintroduce the
582 ¹H,¹³C, or ¹H,¹⁵N dipolar couplings, e.g., in DIPSHIFT
583 experiments.^{172–175} CPPI type experiments are insensitive
584 with respect to remote protons and yield improved performance
585 in comparison to the noncompensated sequences.^{176–178} CPPI
586 is, however, affected by RF inhomogeneities on the RF channel
587 that is employed to switch between the matching conditions.¹⁷⁹
588 Alternatively, REDOR can be employed instead of CPPI, if
589 remote protons can be depleted as in perdeuterated
590 samples.^{179–181} A further possibility to determine the motional
591 amplitude at a given site are off-magic-angle spinning experi-
592 ments with angle offsets as small as 0.03°. ^{182,183} At 100 kHz
593 MAS, motionally averaged dipolar couplings can be determined
594 using variable-contact CP.¹⁸⁴

To quantitatively describe dynamics, the extended model-free
596 formalism has been introduced.¹⁸⁵ At least four observables are
597 needed to yield a quantitative description of motion on two time
598 scales and to fit motional amplitudes and correlation times. For
599 the fit of the data for α-SH3,⁸² a combination of ¹⁵N-R₁
600 relaxation rates measured at different fields,^{186,187} ¹H–¹⁵N
601 dipole, ¹⁵N-CSA cross-correlated relaxation rates η,^{74,188,189} as
602 well as CPPI derived dipolar order parameters¹⁹⁰ were
603 employed. Spin diffusion can potentially affect the quantification
604

605 of R_1 relaxation rates.^{187,191,192} If magnetization flows to remote
606 nuclei, e.g., to methyl groups which have a short relaxation time,
607 the relaxation rate will appear systematically larger. Cross-
608 correlated relaxation rates η yield frequency independent
609 spectral density functions, similar to transversal R_2 rates in
610 solution. As an alternative, $R_{1\rho}$ rates can be analyzed to yield site-
611 specific incoherent R_2 type relaxation rates.¹⁹³ Depending on the
612 resonance offset, $R_{1\rho}$ depends on a combination of R_1 , R_2 , and
613 the exchange rate in case slow motional processes are involved.
614 In principle, the number of observables can be increased if
615 relaxation parameters measured for neighboring nuclei are taken
616 into account. Lamley et al. have shown that the motion of the
617 peptide plane can be modeled from combined backbone ^{15}N
618 amide and $^{13}\text{C}'$ carbonyl data.¹⁹⁴ This allows increase of the
619 number of time scales in the extended model free analysis. Using
620 this extended model-free formalism,¹⁸⁵ the dynamics of the α -
621 spectrin SH3 domain,⁸² ubiquitin,^{195,196} and superoxide
622 dismutase¹⁵³ were quantified. Similarly, backbone dynamics
623 was characterized for the D76N mutant of β -2 microglobulin.¹⁹⁷
624 R_1 , $R_{1\rho}$, and order parameters were employed to quantify
625 dynamics in an extended model-free analysis. It was shown that
626 the outer strands of D76N β 2m are destabilized accounting for
627 the increased aggregation propensity of this mutant. Site-specific
628 R_1 , $R_{1\rho}$, and order parameters via REDOR have been measured
629 further for the backbone ^{15}N and $^{13}\text{C}^\alpha$ nuclei for fibrils of the
630 fungal prion protein HET-s(218–289).¹⁹⁸ The ^{15}N and $^{13}\text{C}^\alpha$
631 data were fit, assuming motions at three time scales which
632 provides a statistically significant, better fit of the relaxation data.
633 There, a slow correlation time of 6.2 μs for ^1H – ^{15}N , and 4.1 μs
634 for $^{13}\text{C}^\alpha$ – $^1\text{H}^\alpha$ – $^{13}\text{C}^\alpha$ are obtained, indicating the presence of a
635 slow global fibril motion. The associated order parameters are
636 very high, with $S_S^2 \geq 0.98$.

637 The dilemma with the analysis of relaxation rates in the solid
638 state lies in the problem that R_1 spin–lattice relaxation rates are
639 determined not only by incoherent effects, i.e., structural
640 fluctuations, but to a large extent also by coherent interactions
641 that are dependent on the local spin density. Care has to be taken
642 to disentangle the involved coherent and incoherent effects.
643 Isolated nuclei are in principle ideally suited to quantify
644 dynamics because inhomogeneous interactions are more easily
645 refocused. Spin diffusion, e.g., during a relaxation period, results
646 in homonuclear magnetization transfer and yields an apparently
647 shorter ^{13}C T_1 relaxation time in case a nucleus is interacting
648 with, e.g., a methyl group which has a short relaxation time and
649 basically acts as a “sink” for magnetization.¹⁹⁹ For a correct
650 analysis of backbone C^α carbon relaxation rates, dilution of both
651 the proton and carbon spin system together with fast MAS are
652 thus important.²⁰⁰ A dilution of the carbon spin system can be
653 achieved by growing the protein in a M9 minimal medium that
654 employs [μ - ^2H , 2- ^{13}C]- or [μ - ^2H , 1,3- ^{13}C]-glycerol as the sole
655 carbon source. To introduce protons sporadically, the medium
656 contains low amounts of H_2O .³⁹ A set of MAS dependent
657 relaxation measurements show that both proton and carbon spin
658 dilution are necessary to yield a MAS independent and
659 monoexponential, e.g., ^{13}C - T_1 relaxation decay curves. Alter-
660 natively, selective methyl labeling allows to incorporate an
661 isolated carbon spin into the protein.³²

662 In the Lipari–Szabo model-free approach, motion is
663 described by employing a sum of spectral density functions.
664 Each term of the sum reflects a linear independent motional
665 regime and contains an order parameter and a correlation time.
666 Smith et al.²⁰¹ hypothesize that the fit might be distorted, and
667 specifically those correlation times are obtained where relaxation

experiments are particularly sensitive. The model-free approach 668
requires linear independent motional regimes. If dynamics are 669
more complex, the model-free approach breaks down. This 670
induces a bias in the resulting dynamics description. It was 671
therefore suggested to employ so-called “dynamics detectors”, 672
which characterize different ranges of correlation times. In this 673
approach, no specific model for the correlation function is 674
needed any longer to analyze protein motion. 675

Rotating-frame relaxation rates ($R_{1\rho}$) rates can be exploited 676
further to probe micro- to millisecond time scale mo- 677
tions.^{202–204} $R_{1\rho}$ measurements have been employed to analyze 678
protein rocking motions in crystals.^{205–208} At low RF field 679
strengths in the $R_{1\rho}$ experiment, isotropic chemical shift 680
fluctuation induce the relaxation dispersion that are familiar 681
from solution-state NMR. When the effective RF field ω_e 682
approaches the rotary resonance condition ($\omega_e \simeq n \times \omega_r$, with 683
 $n = 1, 2$) bond angle fluctuations induce an increased $R_{1\rho}$ rate 684
reflecting fluctuations of anisotropic interactions. This condition 685
is referred to as near-rotary-resonance $R_{1\rho}$ relaxation-dispersion 686
(NERRD) and can be exploited in a complementary way to 687
retrieve information on the chemical shift differences between 688
the two states, the amplitude of the motion, the populations, and 689
the involved exchange rates.^{207,209,210} Using this method, the 690
relaxation dispersion profiles for both the Bloch–McConnell 691
and the near-rotary resonance regimes were modeled for fibrils 692
formed by the human prion protein Y145Stop. Under the 693
assumption of a two-state exchange mechanism, the amino acids 694
of the fibril core exchange with a rate on the order of 1000 s^{-1} 695
with a population of the excited state of 2%.²¹¹ Similarly to ^{15}N 696
 $R_{1\rho}$, also ^1H $R_{1\rho}$ relaxation-dispersion can be measured for highly 697
deuterated protein samples.^{212,213} While ^{15}N $R_{1\rho}$ relaxation rates 698
rather sense local fluctuations, ^1H $R_{1\rho}$ is a sensor of more global 699
structural rearrangements due to its higher gyromagnetic ratio. 700
Using these experiments, loop motions have been observed, 701
suggesting that the protein adopts a binding-competent 702
conformation in dynamic equilibrium with a sterically impaired 703
ground-state conformation both in solution and in its crystalline 704
form. 705

Additional information on the underlying motional model can 706
come from heteronuclear NOE experiments. A significant 707
heteronuclear NOE is expected for side chain methyl and amino 708
groups due their fast 3-fold rotation which results in a fluctuation 709
of the effective dipolar coupling at the respective heteronucleus 710
and thus in relaxation. The heteronuclear NOE in combination 711
with cross-polarization can be exploited for polarization 712
transfer.²¹⁴ This has subsequently been applied to prote- 713
teins.^{215–217} Heteronuclear NOEs for amide nitrogens in the 714
protein backbone in the solid state were observed first by Giraud 715
et al.²¹⁸ Site-specific [^1H] ^{13}C and [^1H] ^{15}N heteronuclear NOE 716
rates for methyl and amide backbone as well as for side chain 717
groups were quantified by Lopez del Amo and co-workers.²¹⁹ 718
The heteronuclear NOE is sensitive only to very fast time scale 719
motion and enables the detection of motional processes with a 720
picosecond correlation time. The low sensitivity of the 721
experiment, however, prevents a more detailed quantitative 722
analysis. More recently, the heteronuclear NOE involving 723
methyls has been exploited in high-field DNP experiments²²⁰ 724
as a dynamic sensors for probing local molecular packing and to 725
selectively enhance the polarization of nuclei of residues within 726
the ligand-binding pocket. 727

MD simulations can help to visualize motion and to get 728
quantitative motional models. For the α -spectrin SH3 domain-, 729
solution-, and solid-state NMR derived order parameters were 730

731 quantified and compared to a MD trajectory to quantify the
732 amplitude and time scale of backbone motion.²²¹ It was found
733 that in α -spectrin SH3 ns- μ s motions occur in dynamic loops,
734 termini, and side chains, but generally not in the structured
735 portion of the backbone. Similarly, combination of relaxation
736 data and MD simulations allowed to characterize the dynamics
737 of fibrils of the fungal prion protein HET-s(218–289).²²² Here,
738 the NMR data were subjected to the dynamics detector method.
739 From the MD trajectory, cross-correlation coefficients were
740 extracted to find out whether motion of different residues occurs
741 on the same time scale in a correlated way. It was found that local
742 motions occur typically on short time scales, while for longer-
743 range correlated motion take place on longer time scales. Grohe
744 et al. compared MD simulations, NMR relaxation measure-
745 ments, and an eNOE-based structural ensemble approach for
746 the chicken α -spectrin SH3 domain. The combined analysis
747 showed that the data are consistent with respect to both the time
748 scales and the conformational states that are sampled in the
749 dynamic MD trajectory.²²³

750 In the following, a few examples shall be given in which efforts
751 have been made to quantify amplitude and time scale for
752 different biomolecular systems. For ubiquitin, 1 μ s time scale
753 motion has been observed for the entire protein. Most residues
754 displayed small-amplitude motion, while larger-amplitude
755 dynamics were observed for the β 1- β 2 turn and the N terminus
756 of the α -helix.²²⁴ For the β -barrel OmpA, ^1H - ^{15}N dipolar
757 coupling as well as ^{15}N R_1 and $R_{1\rho}$ relaxation rates at fast (60
758 kHz) MAS and high magnetic field (1 GHz) were measured to
759 describe the motion site specifically.²²⁵ It has been suggested
760 that motion occurs as a collective rocking motion of low
761 amplitude and of hundreds of nanoseconds time scale.
762 Enhanced ^{15}N $R_{1\rho}$ relaxation rates observed for residues located
763 at the edges of the strands suggest an increased mobility of the
764 connecting loops. By determining the de novo NMR structure of
765 the membrane protein AlkL in lipid bilayers and by accessing
766 structural dynamics in combination with MD simulations, MAS
767 solid-state NMR experiments allowed to solve the controversy
768 between crystal and NMR structures of homologues of this
769 protein.^{226,227} It was found that dynamic lateral exit sites occur
770 through restructuring of a barrel extension formed by the
771 extracellular loops. Investigations of the fully mature functional
772 Cu, Zn state of superoxide dismutase and its E, Zn-state in which
773 the Cu site is empty have shown that the metal ion after uptake
774 does not act as a rigidification element. Using ^{15}N R_1 , $R_{1\rho}$
775 experiments, it has been shown that metal binding acts rather as
776 a switch that redistributes motional processes on different time
777 scales and induces a coupling of the dynamics of the histidine
778 side chains to remote key backbone elements of the protein.²²⁸
779 The SAIL approach allows selectivity of isotopically labeled
780 aromatic rings in proteins.¹⁶⁰ This labeling in combination with
781 ^1H detected experiments yielded the quantification of motional
782 amplitudes and rates of phenylalanine side chain dynamics in the
783 half-megadalton enzyme complex, aminopeptidase TET2.¹⁶¹ In
784 the past, carbonic anhydrase II was considered to be a
785 prototypical rigid drug target. While X-ray crystallography has
786 suggested that CA-II has little structural plasticity, MAS solid-
787 state NMR could demonstrate that the binding pocket
788 undergoes pronounced open/close conformational-exchange
789 dynamics, which are abrogated upon binding of a sulfonamide
790 inhibitor.^{229,230} ^1H detection is further employed in ligand
791 titration experiments to find out how the ClpP protease
792 machinery is activated by substrates or blocked by active-site
793 inhibitors.²³¹

11. MOTIONAL INTERFERENCE WITH DECOUPLING

The investigations of protein dynamics in the solid state
described above are not purely academic but are of fundamental
practical importance. Local dynamics in the ns- μ s regime with a
significant amplitude will have detrimental effects on the quality
of the obtained correlation spectra recorded with high-power
decoupling. Averaging of $\text{H}^\alpha/\text{H}^\beta$ spin states results then in a
broadening of the observed, e.g., ^{15}N resonances if the respective
amide moiety undergoes large amplitude ns- μ s motion.
Similarly, proton resonances also would be broadened in case
the heteronuclear spin is decoupled.²³² This effect was observed
for several residues in the α -spectrin SH3 domain, which have a
local correlation time on the order of several hundred
nanoseconds.^{74,233} For these residues, it turns out that
TROSY²³⁴ like experiments yield spectroscopic advantages. In
fact, broad proton lines were observed previously for fibrils of
perdeuterated Alzheimer's disease $\text{A}\beta$ 40.^{52,235} For fibrils of α -
synuclein fibrils²⁶ and hydrophobins,²³⁶ a similar spectral quality
is obtained. Spin-state selective excitation and transfer experi-
ments in principle allow to observe the narrow multiplet
component. However, so far, no beneficial effects from TROSY
type experiments have been observed for amyloid fibrils or
membrane proteins in the solid state.⁸⁷

12. SEDIMENTED SOLUTIONS

In addition to classical biological solids such as amyloid fibrils
and membrane proteins, MAS solid-state NMR allows to
investigate high molecular weight protein complexes in solution.
To achieve refocusing of anisotropic interactions by magic angle
spinning, the tumbling correlation time of the protein in the
rotor must exceed the rotor period.²³⁷ In a typical MAS solid-
state NMR experiment, rotation of the sample induces
accelerations due to centrifugation on the order of 1–10 \times
10⁶ g.^{238,239} Depending on its molecular weight, sedimentation
efficiently immobilizes the soluble protein. Essential for the
success of the method is the development of packing tools that
allow efficient transfer of protein material into the solid-state
NMR rotor.^{240–245} Using this approach, the 20S proteasome,²⁴⁶
the 50S ribosome,^{247,248} the 50S bound trigger factor,²⁴⁹ and
immunoglobulins in complex with GB1⁶⁰ have been inves-
tigated. Small heat shock proteins such as α B-crystallin assemble
into oligomers with a molecular weight of 400–600 kDa. These
complexes are too heterogeneous and dynamic to form stable
crystals that can be analyzed with X-ray crystallography.²⁵⁰
Misfolding peptides and proteins interact with the chaperones
and form chaperone-substrate complexes in the MDa range.
These substrate complexes remain in solution, in contrast to the
pure substrates that aggregate and precipitate. Traditionally,
protein complexes are crystallized using a precipitant such as
polyethylene glycol (PEG).²⁵¹ This procedure prevents the
investigation of chaperone-substrate complexes as the inter-
action between the small heat shock protein and the
amyloidogenic substrate is very weak. MAS solid-state NMR is
as such the perfect tool to investigate chaperone-amyloid
complexes: MAS solid-state NMR allows overcoming of the
limitations of NMR imposed by the tumbling of the proteins in
solution. At the same time, sedimentation is not irreversible, but
samples can be recovered after the experiment. The inves-
tigations can be carried out in a native-like environment without
the need to irreversibly precipitate the sample.

Immobilization of protein complexes by sedimentation
enabled furthermore the investigation of transient complexes

853 formed by the 177 residue homotetrameric single-stranded
854 DNA binding protein SSB with the single-stranded DNA that is
855 involved in regulation of *Escherichia coli* DNA metabolism.²⁵²
856 Using proton-detected solid-state NMR experiments, the
857 binding site of the inhibitor bortezomib could be identified at
858 an active-site residue on the 14-subunit complex of caseinolytic
859 protease ClpP.²³¹ Finally, the method allows to study
860 membraneless cellular organelles, liquid–liquid phase separa-
861 tion, and hydrogels.²⁵³

13. CONSEQUENCES OF DEUTERATION ON PROTEIN STABILITY, ENZYMATIC ACTIVITY, AND FUNCTION

864 Deuteration can in principle affect the chemical properties of a
865 protein. It has been shown that D₂O and thus replacing protons
866 by deuterons at exchangeable sites has a stabilizing effect on the
867 protein structure.^{254–256} On the other hand, substitution of
868 protons with deuterons at nonexchangeable sites decreases the
869 hydrophobic effect,^{257,258} and this way potentially increases
870 dynamics within the hydrophobic core of a protein. For
871 transthyretin, it has been shown that deuteration influences
872 subunit exchange and accelerates its aggregation kinetics.²⁵⁹
873 Similarly, the hydrophobic effect can affect the activity of an
874 enzyme.²⁶⁰ Other reports have shown, however, that enzyme
875 activity is not affected by deuteration.^{261,262} In any case, care
876 should be taken and the macroscopic properties of protonated
877 and deuterated samples have to be compared in order to rule out
878 systematic errors in the analysis.

14. FASTER MAS, HIGHER FIELDS, AND PROTONATED SAMPLES

880 Initially, proton detected experiments were carried out using
881 perdeuterated samples in which 20% of the exchangeable sites
882 were protonated.^{25,79} The employed 3.2 mm rotors achieve
883 rotation frequencies on the order of 10–20 kHz and can
884 accommodate 40 μ L of sample. Faster MAS has historically led
885 to higher quality and more informative solid-state NMR spectra.
886 An increase of the MAS rotation frequency to 60 kHz was
887 sufficient to increase the proton concentration so that 100% of
888 all exchangeable sites could be protonated without loss of
889 resolution in the proton dimension.^{78,263} The 1.3 mm rotor,
890 which is necessary to achieve a rotation frequency of 60 kHz, can
891 only be filled with 4 μ L of material and thus 10 \times less sample. The
892 anticipated loss in sensitivity is more than compensated with the
893 5 \times increase of the labeling density at exchangeable sites and the
894 2.5 \times higher quality factor of the NMR coil, which increases with
895 the inverse of the diameter. The effective sensitivity of the 1.3
896 mm rotor is thus actually 25% larger than the sensitivity achieved
897 by a 3.2 mm rotor despite the fact that only 10% of the material is
898 used. Using a 0.7 mm MAS rotor and rotation frequencies
899 around 110 kHz enable the analysis of a fully protonated
900 sample.³ The active volume of a 0.7 mm rotor amounts to ca. 0.4
901 μ L. Under these conditions, the absolute intensity of an amide
902 proton resonance from a protonated protein will be lower than
903 the respective intensity from a perdeuterated samples in a 1.3
904 mm rotor, as the proton concentration at a particular site cannot
905 be increased further. However, the 0.7 mm coil has a 2 \times
906 improved coil efficiency. At the same time, T_2' and $T_{1\rho}$ increase
907 in smaller rotors,¹¹¹ which compensates for some of the losses
908 implicated by the smaller rotor volumes.

909 The maximum MAS frequency is achieved if the rotor surface
910 approaches the speed of sound.²⁶⁴ Thus, higher MAS

frequencies can only be achieved if the rotor diameter decreases. 911
In the future, faster MAS can potentially be achieved by making 912
use of rotors that have a different geometry²⁶⁵ or by employing 913
gases that are less viscous than nitrogen. However, the latter 914
point remains to be shown because the strength of the ceramic 915
material has to increase together with the rotation frequency. 916
Ultrafast MAS experiments are potentially interesting for 917
applications, where it is difficult to produce large amount of 918
labeled and reconstituted protein in case the overall sensitivity 919
does not suffer. Faster rotation yields better averaging of proton 920
dipolar interactions. Recently, an experimental MAS frequency 921
dependent intensity analysis for the fully protonated proteins 922
ubiquitin, Rpo4/7, HET-s(218–289), and Cp149 was carried 923
out.^{266,267} Comparison of proton solid-state NMR spectra 924
recorded at 100 and 150 kHz MAS shows that the average 925
proton homogeneous line width is reduced by a factor of 1.5, 926
while the experimental line width is decreased by a factor 927
1.25.²⁶⁷ At the same time, the sensitivity of an HN correlation 928
experiment is increased by a factor of 1.48. This factor accounts 929
for differences in the efficiencies of the polarization transfer 930
steps, the sample volumes, and coil efficiencies of the two probes 931
as well as line narrowing by MAS. 932

Direct observation of protons and development of fast 933
spinning probes raises the question which maximum MAS 934
frequency is actually required to achieve the optimum sensitivity 935
for a particular (protonated) sample. For a highly deuterated 936
protein sample in which only methyl groups are selectively 937
protonated, the MAS rotation frequency that is required to yield 938
half of the maximum possible sensitivity ranges from 20 kHz up 939
to 324 kHz, depending on the specific site with an average value 940
of (135 \pm 88) kHz at a magnetic field strength of 11.7 T (500 941
MHz).³⁸ In case a sensitivity is supposed to be reached that 942
corresponds to 80% of the maximum possible sensitivity, a MAS 943
rotation frequency of (498 \pm 370) kHz is necessary. It is easy to 944
imagine that much higher rotation frequencies are necessary to 945
achieve the same intensity for a protonated protein sample. 946

To appreciate the effect of the proton density on the proton 947
spectral quality, microcrystalline perdeuterated protein samples 948
have been investigated in which protons were selectively 949
introduced into methyl groups.³⁸ Specifically, methyl groups 950
were labeled either as CH₃, CH₂D, or CHD₂ isotopomers. This 951
allows to modulate the proton density in the protein. 952
Subsequently, the MAS frequency dependent cross peak 953
intensities of the experimental and theoretical spectra was 954
analyzed. It was found that the break-even point, i.e., the MAS 955
frequency that is necessary to reach equal intensity in CH₃ 956
versus CHD₂ selectively methyl protonated samples is depend- 957
ent on the specific methyl group and is achieved on average at a 958
MAS frequency of (75 \pm 53) kHz assuming a B₀ field of 1 GHz. 959
Similarly, on average, a MAS frequency of (41 \pm 28) kHz is 960
necessary to achieve the same intensity for CH₂D versus CHD₂ 961
selectively methyl protonated samples. This information is 962
potentially of interest for experiments that involve very large, 963
nonsymmetric protein complexes where signal-to-noise is 964
critical and for which any additional increase in sensitivity is 965
essential for the success of the experiment.^{60,247–249} 966

In the past, more powerful magnets and higher frequency 967
MAS has given access to new areas of research. Novel 968
superconducting magnet technology^{268,269} made magnetic fields 969
available in the range of 1.2 to 1.5 GHz. Spectral sensitivity scales 970
with B₀^{3/2}.²⁷⁰ For proton resonances in solid proteins, sensitivity 971
gains beyond the B₀^{3/2} factor are expected when the difference of 972
the isotropic chemical shift is larger than the size of the ¹H–¹H 973

974 dipolar coupling. In a showcase study, this has been
975 demonstrated for a selectively methyl protonated α -SH3
976 domain sample in a deuterated background.²⁷¹ It is found that
977 for residues which are embedded in a dense proton coupling
978 network, sensitivity is 2-fold larger than what is expected from
979 the canonical field dependence. The gains are largest for methyl
980 groups that experience large proton dipolar couplings and are
981 thus not dependent on instrumental issues that determine the
982 efficiency of the spectrometer, of the probe, and of the
983 polarization transfer. Recent experimental results at 28.2 T
984 (1.2 GHz) suggest that in particular proton detected MAS solid-
985 state NMR experiments benefit dramatically from the increased
986 magnetic field strength.^{272,273}

987 At 110 kHz MAS, deuterated samples (including selectively
988 methyl protonated samples) yield better sensitivity and
989 resolution in comparison to protonated proteins.²⁷⁴ It remains
990 to be seen which further improvements the recent developments
991 in magnet and probe design will have on the proton spectral
992 quality in MAS solid-state NMR experiments.

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1004 Notes

1005 The author declares no competing financial interest.

1006 Biography

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1011 Institute of Technology, Cambridge, USA, he was heading an Emmy-
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1016 Charité Universitätsmedizin Berlin. In 2010, Dr. Reif accepted an offer
1017 from the Department of Chemistry at the Technische Universität
1018 München and became coaffiliated at the Helmholtz-Zentrum München
1019 in Neuherberg. His research interests involve the development of MAS
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1021 dynamical properties of proteins in the solid-state, as well as the
1022 application of solution- and solid-state NMR spectroscopy to the study
1023 of amyloidogenic peptides and proteins.

1024 ACKNOWLEDGMENTS

1025 B.R. acknowledges support from the Helmholtz-Gemeinschaft,
1026 the Deutsche Forschungsgemeinschaft (DFG, grant Re1435).

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